

# Quantum Mechanics 2

*Robert C. Roleda*  
*Physics Department*

Hydrogen Atom  
Zeeman Effect



De La Salle University



# Normal Zeeman Effect

If a Hydrogen atom is placed in a magnetic field, it interacts with this field through its orbital and spin angular momenta.

$$H_Z = -\mu \cdot B$$

The magnetic moment of an electron due to its orbital angular momentum is

$$\mu_l = -\frac{e}{2m_e} L$$

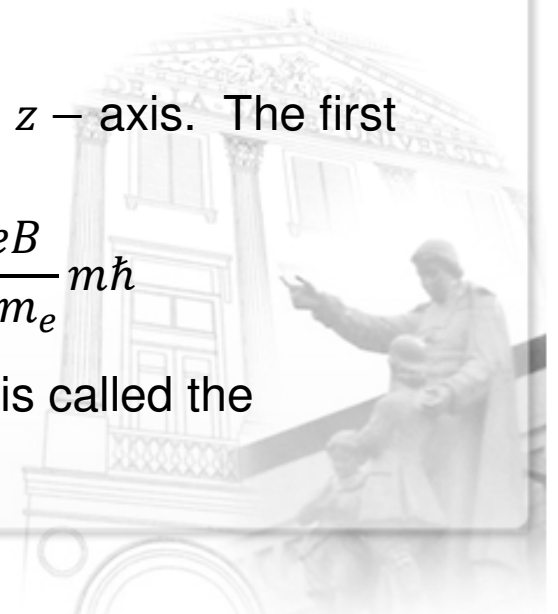
Thus,

$$H_Z = \frac{e}{2m_e} L \cdot B$$

For a uniform magnetic field, we may take its direction as the  $z$  – axis. The first order correction due to this interaction is then

$$E_{(NZ)}^{(1)} = \langle nlm | H_Z | nlm \rangle = \frac{eB}{2m_e} \langle nlm | L_Z | nlm \rangle = \frac{eB}{2m_e} m\hbar$$

This leads to the splitting of the azimuthal degeneracy, and it is called the **Normal Zeeman Effect**



# Anomalous Zeeman Effect

The magnetic moment of an electron due to its spin is

$$\mu_s = -\frac{e}{m_e} S$$

Its interaction Hamiltonian is then

$$H_Z = \frac{e}{m_e} S \cdot B$$

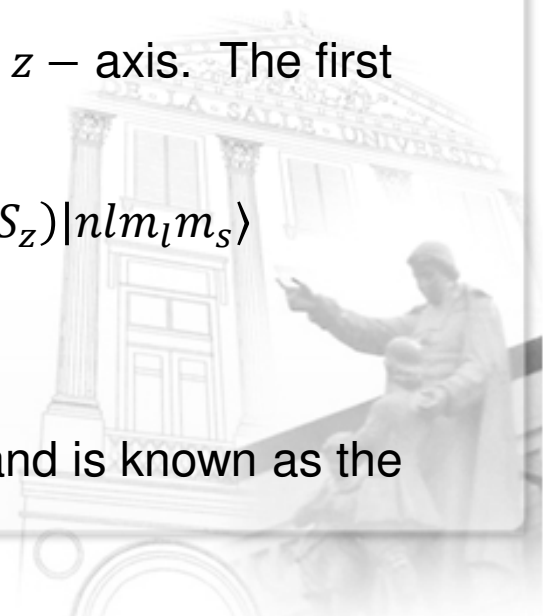
Taken together with the interaction with the orbital angular momentum

$$H_Z = \frac{e}{2m_e} (L + 2S) \cdot B$$

For a uniform magnetic field, we may take its direction as the  $z$  – axis. The first order correction due to this interaction is then

$$\begin{aligned} E_{(AZ)}^{(1)} &= \langle nlm_l m_s | H_Z | nlm_l m_s \rangle = \frac{eB}{2m_e} \langle nlm_l m_s | (L_z + 2S_z) | nlm_l m_s \rangle \\ &= \frac{eB}{2m_e} (m_l + 2m_s) \hbar \end{aligned}$$

This leads to a doubling of the azimuthal degeneracy splits, and is known as the **Anomalous Zeeman Effect**



# Weak Fields

If the external magnetic field is very weak, the Zeeman interaction is much weaker than the spin-orbit interaction. Thus, we may take it as a perturbation relative to the Hamiltonian

$$H_0 = \frac{p^2}{2m_e} - \frac{e^2}{4\pi\epsilon_0 r} + \frac{e^2}{4\pi\epsilon_0} \frac{1}{2m_e^2 c^2 r^3} S \cdot L$$

The good quantum numbers for fine structure are  $n, l, s, j, m$ . Thus,

$$E_{(WZ)}^{(1)} = \langle nlsjm | H_Z | nlsjm \rangle = \frac{eB}{2m_e} \langle nlsjm | (L_z + 2S_z) | nlsjm \rangle$$

The total angular momentum is

$$J = L + S$$

Thus,

$$L_z + 2S_z = J_z + S_z$$

The first term is compatible with the quantum numbers but the second is not



# Spin

We note that the total angular momentum is a constant of motion, and that  $L$  and  $S$  have fixed magnitudes but the vectors precess about the direction of  $J$ .

We may then express

$$\vec{S} = (\vec{S} \cdot \hat{J})\hat{J} = \frac{(\vec{S} \cdot \vec{J})}{J^2} \vec{J}$$

Now,

$$L = J - S$$

so

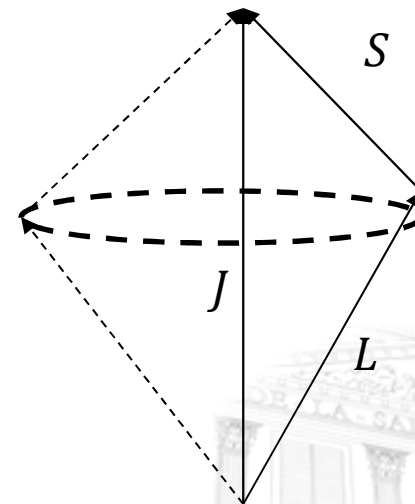
$$L^2 = J^2 + S^2 - 2S \cdot J$$

Thus,

$$S \cdot J = \frac{1}{2} [J^2 + S^2 - L^2]$$

and

$$S_z = \frac{[J^2 + S^2 - L^2]}{2J^2} J_z$$



# Weak Field Correction

Therefore, in weak fields

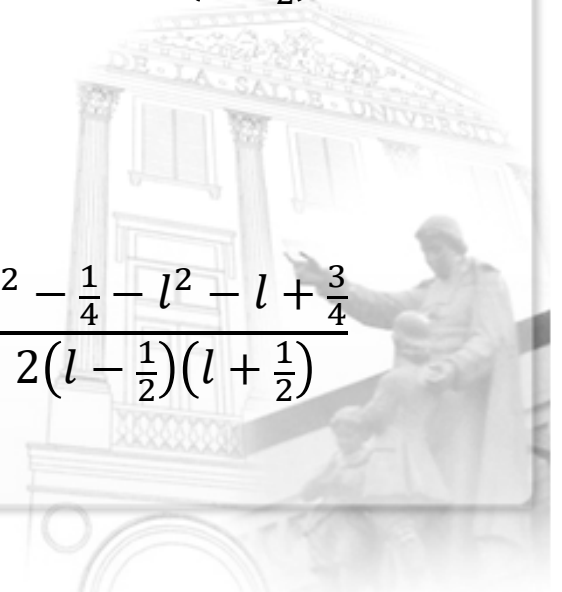
$$\begin{aligned} E_{(WZ)}^{(1)} &= \frac{eB}{2m_e} \left\langle nlsjm \left| \left( J_z + \frac{[J^2 + S^2 - L^2]}{2J^2} J_z \right) \right| nlsjm \right\rangle \\ &= \frac{eB}{2m_e} \left[ 1 + \frac{j(j+1) + \frac{3}{4} - l(l+1)}{2j(j+1)} \right] m\hbar \end{aligned}$$

For  $j = l + \frac{1}{2}$ ,

$$\begin{aligned} \frac{j(j+1) + \frac{3}{4} - l(l+1)}{2j(j+1)} &= \frac{(l + \frac{1}{2})(l + \frac{3}{2}) - l(l+1) + \frac{3}{4}}{2(l + \frac{1}{2})(l + \frac{3}{2})} = \frac{l^2 + 2l + \frac{3}{4} - l^2 - l + \frac{3}{4}}{(2l+1)(l + \frac{3}{2})} \\ &= \frac{l + \frac{3}{2}}{(2l+1)(l + \frac{3}{2})} = \frac{1}{(2l+1)} \end{aligned}$$

For  $j = l - \frac{1}{2}$ ,

$$\begin{aligned} \frac{j(j+1) + \frac{3}{4} - l(l+1)}{2j(j+1)} &= \frac{(l - \frac{1}{2})(l + \frac{1}{2}) - l(l+1) + \frac{3}{4}}{2(l - \frac{1}{2})(l + \frac{1}{2})} = \frac{l^2 - \frac{1}{4} - l^2 - l + \frac{3}{4}}{2(l - \frac{1}{2})(l + \frac{1}{2})} \\ &= \frac{-l + \frac{1}{2}}{(l - \frac{1}{2})(2l+1)} = -\frac{1}{(2l+1)} \end{aligned}$$



# Weak Field Correction

The weak field Zeeman effect is therefore

$$E_{(WZ)}^{(1)} = \frac{eB}{2m_e} \left[ 1 \pm \frac{1}{(2l + 1)} \right] m\hbar$$

